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E. V. Laitone [2] has called our attention to a useful (apparently not so well advertised) procedure for obtaining an analytic function from one of its conjugate harmonics. The idea is simply this: if $f(z) = f(x + iy) = \phi(x, y) + i\psi(x, y)$ is analytic in a neighborhood of some point $z_0 = x_0 + iy_0$, then it can be recovered from either of its conjugate harmonics through the integration formulas,

$$f(z) = \int \left[\frac{\partial \phi}{\partial x} \left(z - i y_0, y_0 \right) - i \frac{\partial \phi}{\partial y} \left(z - i y_0, y_0 \right) \right] dz = \int \left[\frac{\partial \psi}{\partial y} \left(z - i y_0, y_0 \right) + i \frac{\partial \psi}{\partial x} \left(z - i y_0, y_0 \right) \right] dz \qquad (1)$$

in the "variable" z. These formulas are readily obtained from the Cauchy-Riemann equations once it is observed that

$$f(z) = \phi(z - iy_0, y_0) + i\psi(z - iy_0, y_0)$$
(2)

holds for all z in some neighborhood of z_0 .

Actually, Laitone derived (2) only for $y_0 = 0$, using the conjugate complex function $\overline{f}(z)$. An alternative (and perhaps more direct and illuminating) derivation can be given using Taylor series. For if

$$\phi(u + x_0, y_0) = \sum a_n u^n, \, \psi(u + x_0, y_0) = \sum b_n u^n \tag{3}$$

are the Taylor series expansions (in the first variable) of ϕ and ψ about x_0 , then

$$a_n = \frac{1}{n!} \frac{\partial^n \phi}{\partial x^n} (x_0, y_0), b_n = \frac{1}{n!} \frac{\partial^n \psi}{\partial x^n} (x_0, y_0)$$

hold for all *n*. But $d^n f/dz^n = \partial^n \phi/\partial x^n + i \partial^n \psi/\partial x^n$, so that the Taylor series expansion of *f* about $z_0 = x_0 + iy_0$ gives

$$f(z) = \sum \frac{1}{n!} \frac{d^n f}{dz^n} (z_0) [z - z_0]^n = \sum (a_n + ib_n) [z - (x_0 + iy_0)]^n = \phi(z - iy_0, y_0) + i\psi(z - iy_0, y_0)$$

where we have used (3) with $u = z - (x_0 + iy_0)$.

. . .

Using analogous Taylor series and $d^n f/dz^n = (-i)^n \partial^n \phi/\partial y^n + (-i)^{n-1} \partial^n \psi/\partial y^n$, we similarly obtain an analogous identity

$$f(z) = \phi(x_0, -iz + ix_0) + i\psi(x_0, -iz + ix_0).$$
(4)

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This alternative derivation might help to explain some of the symbolism used in (1) (the permanence of the Cauchy-Riemann equations under analytic continuation etc.) and so help (along with Laitone's examples) to rescue the idea from near oblivion.

Even more helpful is the Taylor series (in two variables)

$$\phi(h + x_0, k + y_0) = \sum \frac{1}{n!} \left(h \frac{\partial}{\partial x_0} + k \frac{\partial}{\partial y_0} \right)^n \phi(x_0, y_0).$$

For with h = u/2 and k = -iu/2, we obtain the identity

$$2\phi\left(\frac{u}{2}+x_0,\frac{-iu}{2}+y_0\right)=\sum \frac{2}{n!}\left(\frac{u}{2}\right)^n\left(\frac{\partial}{\partial x_0}-i\frac{\partial}{\partial y_0}\right)^n\phi(x_0,y_0).$$
(5)

But from the Cauchy-Riemann equations we have

$$(\partial/\partial x_0 - i\partial/\partial y_0)\phi(x_0, y_0) = \partial\phi(x_0, y_0)/\partial x_0 + i\partial\psi(x_0, y_0)/\partial x_0 = df(z_0)/dz_0.$$

Hence $(\partial/\partial x_0 - i\partial/\partial y_0)^2 \phi(x_0, y_0) = (\partial/\partial x_0 - i\partial/\partial y_0) df(z_0)/dz_0 = 2d^2 f(z_0)/dz_0^2$, and so by induction, $(\partial/\partial x_0 - i\partial/\partial y_0)^n \phi(x_0, y_0) = 2^{n-1} d^n f(z_0)/dz_0^n$ holds for all $n \ge 1$. The series in (5) then becomes

$$\phi(x_0, y_0) - i\psi(x_0, y_0) + \sum \frac{1}{n!} (d^n f(z_0)/dz_0^n) u^n$$

and with $u = z - z_0$ we obtain what may be the simplest of all recovery formulas,

$$f(z) = 2\phi \left(\frac{z - z_0}{2} + x_0, \frac{-i(z - z_0)}{2} + y_0\right) - \phi(x_0, y_0) + i\psi(x_0, y_0).$$

$$\left(= 2i\psi \left(\frac{z - z_0}{2} + x_0, \frac{-i(z - z_0)}{2} + y_0\right) - i\psi(x_0, y_0) + \phi(x_0, y_0) \right)$$
(6)

The simplicity in (6), of course, lies in the fact that no integrations are required to obtain f from ϕ (or ψ). Once (6) is anticipated, it can be verified directly. For the left-hand and right-hand members are clearly equal when $z = z_0$, and upon using the Cauchy-Riemann equations, their derivatives can be seen to be equal. In fact, $(z - z_0)/2 + x_0 = (z + \overline{z_0})/2$, $-i(z - z_0)/2 + y_0 = (z - \overline{z_0})/2i$, and

$$\frac{d}{dz} 2\phi \left(\frac{z+\overline{z}_0}{2}, \frac{z-\overline{z}_0}{2i}\right) = \frac{\partial\phi}{\partial x} \left(\frac{z+\overline{z}_0}{2}, \frac{z-\overline{z}_0}{2i}\right) - i \frac{\partial\phi}{\partial y} \left(\frac{z+\overline{z}_0}{2}, \frac{z-\overline{z}_0}{2i}\right)$$
$$= \frac{\partial\phi}{\partial x} \left(\frac{z+\overline{z}_0}{2}, \frac{z-\overline{z}_0}{2i}\right) + i \frac{\partial\psi}{\partial x} \left(\frac{z+\overline{z}_0}{2}, \frac{z-\overline{z}_0}{2i}\right)$$
$$= \frac{df}{dz} (X+iY) \quad \text{with} \quad X = \frac{z+\overline{z}_0}{2}, Y = \frac{z-\overline{z}_0}{2i}$$
$$= \frac{df}{dz} (z).$$

The circumstance that X + iY becomes simply z when $X = (z + \overline{z_0})/2$ and $Y = (z - \overline{z_0})/2i$ plays a crucial role here, of course. Similar verifications of (2) and of (4) can be made.

Eq. (6), for $z_0 = 0$, has been derived (somewhat heuristically) in [1] and [3] and shows that a real analytic function ϕ of two variables is harmonic, say near the origin, *iff*

$$\phi(x, y) = 2 \operatorname{Re} \phi\left(\frac{x + iy}{2}, \frac{y - ix}{2}\right) - \phi(0, 0)$$
(7)

holds there. Of considerable interest also is the companion equation giving the conjugate harmonic function,

$$\psi(x, y) = 2 \operatorname{Im} \phi\left(\frac{x + iy}{2}, \frac{y - ix}{2}\right) + \psi(0, 0).$$
(8)

Thus each stream function ψ can be expressed algebraically (without differentiations or integrations) in terms of the corresponding potential function ϕ , and vice versa.

Perhaps the identities (2), (4), (6) and (8) (along with these elementary derivations) should be better known than they appear to be.

REFERENCES

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- [3] L. M. Milne-Thomson, *Theoretical hydrodynamics*, The Macmillan Company, New York, 4th ed., 1962, 5th ed., 1967