# Summation of a Slowly Convergent Series Arising in Antenna Study 

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Abstract. An equivalent series for the slowly convergent series

$$
\sum_{n=1}^{\infty}\left[\int_{-\pi / 2}^{\pi / 2} \cos ^{\alpha} \theta \cos (n \epsilon \sin \theta)\right]^{2} / n
$$

which arises in antenna theory is obtained. The new form is found to consist of two rapidly convergent series for small $\epsilon$.

A recent study of the electromagnetic radiation from cylindrical structures [1], [3] requires the evaluation of a slowly convergent series $S_{1}=\pi^{2} \sum_{n=0}^{\infty}\left[J_{0}(n \epsilon)\right]^{2} / n$ where $J_{0}(n \epsilon)$ is the zeroth order Bessel function, and $\epsilon$ is a small positive constant. The expression above is a special case of the series

$$
\begin{equation*}
S(\alpha)=\sum_{n=1}^{\infty}\left[p_{n}(\alpha)\right]^{2} / n \tag{1}
\end{equation*}
$$

where

$$
\begin{equation*}
p_{n}(\alpha)=\int_{-\pi / 2}^{\pi / 2} \cos ^{\alpha} \theta \cos (n \epsilon \sin \theta) d \theta, \quad \alpha>-1 \text { and } 0<\epsilon \tag{2}
\end{equation*}
$$

One notes that in the case $\alpha=0$

$$
p_{n}(0)=\int_{-\pi / 2}^{\pi / 2} \cos (n \epsilon \sin \theta) d \theta=\pi J_{0}(n \epsilon)
$$

and therefore $S(0)=S_{1}$. The aim of this brief is to obtain a more rapidly convergent series that is equivalent to Eq. (1).

Substituting Eq. (2) into Eq. (1) and interchanging the order of summation and integration results in

$$
\begin{equation*}
S(\alpha)=\int_{-\pi / 2}^{\pi / 2} \cos ^{\alpha} \theta\left\{\int_{-\pi / 2}^{\pi / 2} \cos ^{\alpha} \theta^{\prime}\left[\sum_{n=1}^{\infty} \cos (n \epsilon \sin \theta) \cos \left(n \epsilon \sin \theta^{\prime}\right) / n\right] d \theta^{\prime}\right\} d \theta \tag{3}
\end{equation*}
$$

It is well known that

$$
\begin{equation*}
\sum_{n=1}^{\infty} \cos (n \epsilon \sin \theta) \cos \left(n \epsilon \sin \theta^{\prime}\right) / n=-\frac{1}{2} \ln 2\left|\cos (\epsilon \sin \theta)-\cos \left(\epsilon \sin \theta^{\prime}\right)\right| \tag{4}
\end{equation*}
$$

Employing the Taylor expansion for $\cos y$, the difference of two cosine functions in the vertical bars can be written as

$$
\begin{equation*}
\cos (\epsilon \sin \theta)-\cos \left(\epsilon \sin \theta^{\prime}\right)=\frac{\epsilon^{2}}{4}\left(\cos 2 \theta-\cos 2 \theta^{\prime}\right)(1-A) \tag{5}
\end{equation*}
$$

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where
(6) $A=\frac{2 \epsilon^{2}}{4!}(x+y)-\frac{2 \epsilon^{4}}{6!}\left(x^{2}+x y+y^{2}\right)+\frac{2 \epsilon^{6}}{8!}\left(x^{3}+x^{2} y+x y^{2}+y^{3}\right)-\cdots$,
(7) $x=\sin ^{2} \theta$ and $y=\sin ^{2} \theta^{\prime}$.

Substitution of Eq. (5) into Eq. (4) leads to

$$
\begin{align*}
& \sum_{n=1}^{\infty} \cos (n \epsilon \sin \theta) \cos \left(n \epsilon \sin \theta^{\prime}\right) / n  \tag{8}\\
& \quad=\ln \frac{2}{\epsilon}-\frac{1}{2} \ln 2\left|\cos 2 \theta-\cos 2 \theta^{\prime}\right|-\frac{1}{2} \ln |1-A|
\end{align*}
$$

If $|A|<1$ we can expand the last term in the Taylor series as follows
(9) $\quad \frac{1}{2} \ln |1-A|=\frac{1}{2} \ln (1-A)=-\frac{1}{2}\left(A+\frac{1}{2} A^{2}+\frac{1}{3} A^{3}+\cdots\right)$.

Inserting Eq. (9) in Eq. (8) and acknowledging that $A$ is defined as in Eq. (6) we obtain

$$
\begin{align*}
& \sum_{n=1}^{\infty} \cos (n \epsilon \sin \theta) \cos \left(n \epsilon \sin \theta^{\prime}\right) / n \\
& \quad=\ln \frac{2}{\epsilon}+\sum_{n=1}^{\infty} \cos 2 n \theta \cos 2 n \theta^{\prime} / n+\frac{\epsilon^{2}}{24}(x+y)  \tag{10}\\
& \quad+\frac{\epsilon^{4}}{2880}\left(x^{2}+6 x y+y^{2}\right)+\frac{\epsilon^{6}}{181440}\left(x^{3}+15 x^{2} y+15 x y^{2}+y^{3}\right)+\cdots .
\end{align*}
$$

It is desired to determine the condition for which the inequality $|A|<1$ is satisfied. The exact form of that condition is not known. However, the upper bound of $A$ can be obtained readily. Since the absolute value of $\sin \theta$ is always less than or equal to unity, we see from Eq. (6)

$$
|A| \leqq \frac{4 \epsilon^{2}}{4!}+\frac{6}{6!} \epsilon^{4}+\frac{8}{8!} \epsilon^{6}+\cdots=\frac{\sinh (\epsilon)}{\epsilon}-1
$$

Consequently, if the condition

$$
(\sinh (\epsilon) / \epsilon)-1<1
$$

or

$$
\begin{equation*}
(\sinh (\epsilon) / \epsilon)<2 \tag{11}
\end{equation*}
$$

is satisfied, then the inequality $|A|<1$ is always true. We acknowledge that (11) is more stringent than we really need.

Substituting Eq. (10) into Eq. (3) and interchanging the order of integration and summation for the series $\sum_{n=1}^{\infty} \cos 2 n \theta \cos 2 n \theta^{\prime} / n$ we arrive at

$$
\begin{align*}
S(\alpha)= & C_{0}{ }^{2} \ln (2 / \epsilon)+f(\alpha)+\frac{\epsilon^{2}}{12} C_{0} C_{1}+\frac{\epsilon^{4}}{1440}\left(C_{0} C_{2}+3 C_{1}^{2}\right) \\
& +\frac{\epsilon^{6}}{90720}\left(C_{0} C_{3}+15 C_{1} C_{2}\right)+\cdots, \tag{12}
\end{align*}
$$

where

$$
\begin{align*}
C_{n} & =\int_{-\pi / 2}^{\pi / 2} \cos ^{\alpha} \theta \sin ^{2 n} \theta d \theta, \quad n=0,1,2, \cdots  \tag{13}\\
f(\alpha) & =\sum_{n=1}^{\infty}\left[\int_{-\pi / 2}^{\pi / 2} \cos ^{\alpha} \theta \cos 2 n \theta d \theta\right]^{2} / n \tag{14}
\end{align*}
$$

Making use of the following definite integrals [2]

$$
\begin{equation*}
\int_{0}^{\pi / 2} \cos ^{\nu-1} z \sin ^{\mu-1} z d z=\frac{1}{2} \frac{\Gamma(\mu / 2) \Gamma(\nu / 2)}{\Gamma((\mu+\nu) / 2)}, \quad \operatorname{Re} \mu>0, \quad \operatorname{Re} \nu>0 \tag{15}
\end{equation*}
$$

and the reflection formula for the Gamma function

$$
\begin{equation*}
\Gamma(z) \Gamma(1-z)=\frac{\pi}{\sin \pi z}, \quad z \neq \text { integer } \tag{16}
\end{equation*}
$$

and Legendre's duplication formula, it can be shown that

$$
\begin{align*}
& C_{0}=\sqrt{ } \pi[\Gamma((1+\alpha) / 2) / \Gamma(1+\alpha / 2)] \\
& C_{n}=\sqrt{ } \pi\left[\Gamma\left(\frac{1+\alpha}{2}\right) / \Gamma\left(1+\frac{\alpha}{2}\right)\right] \prod_{k=0}^{n-1} \frac{2 k+1}{2(n-k)+\alpha}, \quad n=1,2, \cdots . \tag{17}
\end{align*}
$$

Also from [2] we find

$$
\begin{align*}
\int_{-\pi / 2}^{\pi / 2} \cos ^{\alpha} \theta \cos 2 n \theta d \theta &  \tag{18}\\
& =(-1)^{n+1} \frac{2}{\sqrt{ } \pi} \sin \left(\frac{\alpha \pi}{2}\right) \Gamma\left(\frac{1+\alpha}{2}\right) \Gamma\left(1-\frac{\alpha}{2}\right) \frac{d_{n}(\alpha)}{2 n+\alpha}
\end{align*}
$$

where

$$
\begin{align*}
d_{1}(\alpha) & =1 \\
d_{n}(\alpha) & =\prod_{k=1}^{n-1} \frac{2(n-k)-\alpha}{2(n-k)+\alpha}, \quad n \geqq 2, \tag{19}
\end{align*}
$$

and $\alpha>-1$.
Substituting Eq. (18) into Eq. (14) results in

$$
\begin{equation*}
f(\alpha)=\frac{4}{\pi}\left[\sin \frac{\alpha \pi}{2} \Gamma\left(\frac{1+\alpha}{2}\right) \Gamma\left(1-\frac{\alpha}{2}\right)\right]^{2} \sum_{n=1}^{\infty} \frac{\left[d_{n}(\alpha)\right]^{2}}{n(2 n+\alpha)^{2}} . \tag{20}
\end{equation*}
$$

Note that for $-1<\alpha<2, d_{n}(\alpha)$ is a monotonically decreasing function of both $n$ and $\alpha$, and, in particular,

$$
d_{n}(0)=1, \quad d_{n}(1)=\frac{1}{2 n-1}
$$

therefore Eq. (12) is seen to be represented by two rapidly convergent series for small $\epsilon$.

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