# FAST RECURSION FORMULA FOR WEIGHT MULTIPLICITIES ${ }^{1}$ 

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The purpose of this note is to describe and prove a fast recursion formula for computing multiplicities of weights of finite dimensional representations of simple Lie algebras over C.

Until now information about weight multiplicities for all but some special cases [1, 2] has had to be found from the recursion formulas of Freudenthal [3] or Racah [4]. Typically these formulas become too laborious to use for hand computations for ranks $\gtrsim 5$ and dimensions $\gtrsim 100$ and for ranks $\simeq 10$ and dimensions $\simeq 10^{4}$ on a large computer $[5,6]$. With the proposed method the multiplicities can routinely be calculated, even by hand, for dimensions far exceeding these. As an example we present a summary of calculations [7] of all multiplicities in the first sixteen irreducible representations of $E_{8}$.

Let $\mathscr{H}$ be a semisimple Lie algebra over $\mathbf{C}$ with root system $\Delta$ and Weyl group $W$ relative to a Cartan subalgebra $\mathfrak{g}$. Let $\Delta^{+}$be the positive roots with respect to some ordering and $\Pi=\left\{\alpha_{1}, \ldots, \alpha_{l}\right\}$ the set of simple roots. Let $Q$ and $P$ be the root and weight lattices respectively spanning the real vector space $V \subset \wp^{*}$. If $X \subset P$ we denote by $X^{++}$the set of dominant elements of $X$ relative to $\Pi$.

Let $M$ be an irreducible $(\mathscr{H}$-module with highest weight $\Lambda$ and weight system $\Omega$. An important feature of the approach is the direct determination of $\Omega^{++}$without computing outside the dominant chamber. Since every $W$-orbit is represented by one weight $\lambda \in \Omega^{++}$of the same multiplicity, it suffices to compute such $\lambda$ 's.

The recursion formula for computing the multiplicities is a modification (Proposition 4) of the Freudenthal formula in which the Weyl group has been exploited to collapse it as much as possible. After describing the procedure, we present the $E_{8}$ example. Finally the necessary proofs are given.

[^0]Computation of dominant weights and their multiplicities. Examples.
Determination of $\Omega^{++}$. We define inductively a set of disjoint subsets ('layers') of $P^{++}, L_{k}, k=0,1,2, \ldots$, by (1) $L_{0}=\{\Lambda\}, \quad L_{k}=\left\{\gamma \in P^{++}-L_{k-1} \mid \gamma=\lambda-\beta, \lambda \in L_{k-1}, \beta \in \Delta^{+}\right\}$.

Then (Proposition 1) $\bigcup_{k=0}^{\infty} L_{k}=\Omega^{++}$. Thus $\Omega^{++}$can be found directly by computing the layers beginning with $L_{0}$. After $\Omega^{++}$is computed in this way it is reordered according to level. If $\rho^{\nu} \in V$ is defined by $\left(\rho^{2}, \alpha_{i}\right)=1$ for all $i$, then the new (partial) ordering of $\Omega^{++}$is given by the integers $\left(\lambda, \rho^{\vee}\right), \lambda \in \Omega^{++}$.

Computation of the multiplicity $m_{\lambda}$ of $\lambda \in \Omega^{++}$. An $m_{\lambda}$ of level $k$ is given in terms of the multiplicities $m_{\lambda^{\prime}}$ of weights $\lambda^{\prime}$ of levels above the $k$ th one.

Let $\operatorname{Stab}_{W}(\lambda)$ be the stabilizer of $\lambda$ in $W$. Then $\operatorname{Stab}_{W}(\lambda)=W_{T}:=\left\langle r_{i} \mid i \in T\right\rangle$, where $T=\left\{i \mid\left(\lambda, \alpha_{i}\right)=0\right\}$ [8]. Let $\hat{W}_{T}=\left\langle W_{T},-1_{V}\right\rangle$, where $1_{V}$ is the identity transformation on $V . \hat{W}_{T}$ decomposes $\Delta$ into orbits $o_{1}, \ldots, o_{n}$. Each orbit $o_{i}$ contains a unique $\xi_{i}=\Sigma n_{i j} \omega_{j}, \xi_{i} \in \Delta^{+}$and $n_{i j} \geqslant 0$ for all $j \in T$ (Proposition 3). The modified Freudenthal formula is

$$
\begin{equation*}
\sum_{i=1}^{n}\left|o_{i}\right| \sum_{p=1}^{\infty}\left(\lambda+p \xi_{i}, \xi_{i}\right) m_{\lambda+p \xi_{i}}=\left(c_{\Lambda}-c_{\lambda}\right) m_{\lambda} \tag{2}
\end{equation*}
$$

where $\left|o_{i}\right|$ is the number of elements of $o_{i}$ and, for all $\mu \in P$,

$$
\begin{equation*}
c_{\mu}:=(\mu+\rho, \mu+\rho)-(\rho, \rho), \quad \rho:=\frac{1}{2} \sum_{\alpha \in \Delta^{+}} \alpha . \tag{3}
\end{equation*}
$$

The sum on $p$ is in reality finite and by standard properties of weight strings $\lambda+p \xi_{i} \notin \Omega \Rightarrow \lambda+q \xi_{i} \notin \Omega$ for $q>p$.

It is advantageous to work in the $\omega$-basis of the fundamental weights when computing $m_{\lambda}$. Thus writing $\lambda=\Sigma n_{i} \omega_{i}, T=\left\{i \mid n_{i}=0\right\}$. When the positive roots are expressed in this basis one easily determines $\xi_{i}$ 's. With $S_{i}:=$ $\left\{j \in T \mid\left(\xi_{i}, \alpha_{j}\right)=0\right\}$ the orbit sizes $\left|o_{i}\right|$ are given by subgroup indices [ $W_{T}: W_{S_{i}}$ ] or $2\left[W_{T}: W_{S_{i}}\right]$ (Proposition 3).

If in the relation (2) some weight $\mu=\lambda+p \xi_{i}=\Sigma n_{j} \omega_{j}$ is not in $P^{++}$, then some $n_{j}<0$ and $r_{j} \mu=\mu-n_{j} \alpha_{j}$ is on a higher level. A finite number ( $\leqslant\left|\Delta^{+}\right|$) of reflections $r_{i}$ transforms $\mu$ into $\nu \in P^{++}$and $m_{\nu}$ is already computed.

If an extensive computation of weight multiplicities is to be undertaken, it is important to notice that for a given $(\mathscr{H}$ there are only finitely many subsets $T$ of $\{1,2, \ldots, l\}$ and corresponding $\xi_{i}$ and $\left|o_{i}\right|$. It is natural to compute this information once and for all. We are preparing such a table.

Consider an example of the $E_{8}$ representation of dimension 4096000. There are only nine weights in $\Omega^{++}$. In the basis of fundamental weights these are (after reordering according to levels) $\lambda_{0}, \ldots, \lambda_{8}$ (see Table). Here the layers are $L_{0}=\left\{\lambda_{8}\right\}, L_{1}=\left\{\lambda_{7}, \lambda_{6}, \lambda_{3}\right\}, L_{2}=\left\{\lambda_{5}, \lambda_{4}, \lambda_{2}, \lambda_{1}\right\}, L_{3}=\left\{\lambda_{0}\right\}$. Given $\Delta^{+}$in the $\omega$-basis, even by hand, the computation of (4) can be done in a matter of minutes.


Table of weight multiplicities for the first 16 irreducible representations of $E_{8}$ ordered by the levels of their highest weights $\lambda_{0}, \ldots, \lambda_{15}$. Below the step diagonal the $\lambda_{i}$ 's are expressed in terms of the fundamental weights.

The determination of the multiplicity $m_{\lambda_{2}}$ is representative. Suppose we already know $m_{\lambda_{8}}=1, m_{\lambda_{7}}=2, m_{\lambda_{6}}=12, m_{\lambda_{5}}=48, m_{\lambda_{4}}=56, m_{\lambda_{3}}=174$. First we find the quantities $\xi_{i}$ and $\left|o_{i}\right|$ for $\lambda_{2}$. They are $\xi_{1}={ }_{0000001}^{0},\left|o_{1}\right|=28$, $\xi_{2}=-1000000,\left|o_{2}\right|=128, \xi_{3}={ }_{-1000010}^{0},\left|o_{3}\right|=84$. Next we find the weights $\lambda_{2}+p \xi_{i}$ which either are in $\Omega^{++}$or are transformed there by a sequence
of reflections $r_{j}, j=1,2, \ldots, 8$. These weights are $\lambda_{2}+\xi_{1}=\lambda_{6}, \lambda_{2}+\xi_{2}=$ $\lambda_{5}, \lambda_{2}+\xi_{3}=\lambda_{3}$, and $\lambda_{2}+2 \xi_{3}={ }_{-1000020}^{0}$ which after 16 reflections $r_{j}$ is transformed into $\lambda_{7}$. Hence (2) reads

$$
\begin{align*}
\left|o_{1}\right|\left(\lambda_{6}, \xi_{1}\right) m_{\lambda_{6}} & +\left|o_{2}\right|\left(\lambda_{5}, \xi_{2}\right) m_{\lambda_{5}}+\left|o_{3}\right|\left(\lambda_{3}, \xi_{3}\right) m_{\lambda_{3}} \\
& +\left|o_{3}\right|\left(\lambda_{2}+2 \xi_{3}, \xi_{3}\right) m_{\lambda_{7}}=\left(c_{\lambda_{8}}-c_{\lambda_{2}}\right) m_{\lambda_{2}} \tag{4}
\end{align*}
$$

Substituting the corresponding values into (4), one has $28 \cdot 4 \cdot 12+128 \cdot 3 \cdot$ $48+84 \cdot 2 \cdot 174+84 \cdot 4 \cdot 2=(186-96) m_{\lambda_{2}}$ which gives $m_{\lambda_{2}}=552$.

The Table summarizes our results for the 16 irreducible representations of $E_{8}$. A useful check of the results is the equality of dimensions

$$
\begin{equation*}
\operatorname{dim}(M)=\sum_{\lambda_{i} \in \Omega^{++}}\left[W: W_{T_{i}}\right] m_{\lambda_{i}} \tag{5}
\end{equation*}
$$

where the dimensions $\operatorname{dim}(M)$, the number [ $W: W_{T_{i}}$ ] of weights on each $W$-orbit, and the multiplicities $m_{\lambda_{i}}$ are given in the Table.

## Theory.

Proposition 1. Let $M$ be an irreducible $(\stackrel{y}{l}$-module with highest weight $\Lambda$. Then for $\lambda \in P^{++}$with $\lambda \neq \Lambda, \lambda \in \Omega^{++}$if and only if $\left(\lambda+\Delta^{+}\right) \cap \Omega^{++}$ $\neq \varnothing$.

Proof. Suppose that $\lambda \in P^{++}, \alpha \in \Delta^{+}$and $\mu:=\lambda+\alpha \in \Omega^{++}$. For all $\beta \in \Delta^{+},(\lambda, \alpha) \geqslant 0$. Then $(\mu, \alpha)=(\lambda+\alpha, \alpha)>0$. Since the weight string through $\mu$ is $\mu+q \alpha, \ldots, \mu, \ldots, \mu-p \alpha$ where $p-q=2(\mu, \alpha) /(\alpha, \alpha)$, [3], it follows that $p>0$ and hence $\lambda \in \Omega \cap P^{++}=\Omega^{++}$.

Conversely, suppose that $\lambda \in \Omega^{++}, \lambda \neq \Lambda$. We show that there is an $\alpha \in \Delta^{+}$with $\lambda+\alpha \in \Omega^{++}$. There is a $\beta \in \Delta^{+}$with $\lambda+\beta \in \Omega$. If $\lambda+\beta$ is dominant we are done. If not $\left(\lambda+\beta, \alpha_{j}\right)<0$ for some $j$ so by the argument above the $a_{j}$-weight string through $\lambda+\beta$ contains $\lambda+\beta+\alpha_{j}$. Also $\left(\lambda, \alpha_{j}\right) \geqslant 0$ since $\lambda \in P^{++}$so $\left(\beta, \alpha_{j}\right)<0$. Then $\beta+\alpha_{j}$ is a root, $\beta+\alpha_{j} \in \Delta^{+}$, and $\lambda+\beta+$ $\alpha_{j} \in \Omega$. We can replace $\beta$ by $\beta+\alpha_{j}$ in the above and repeat. The process cannot continue indefinitely, so the required $\alpha$ exists.

An interesting consequence of Proposition 1 is
Proposition 2 (Notation of Proposition 1). Let $k$ be the largest integer such that $L_{k} \neq \varnothing$. Then $L_{k}$ is a singleton $\{\omega\}$. Furthermore, $\omega$ depends only on $\Lambda \bmod 0$. In particular $0 \in \Omega$ if and only if $\Lambda \in Q$.

For $T \subset\{1,2, \ldots, l\}$, let $V_{T}:=\Sigma_{i \in T} \mathbf{R} \alpha_{i}$, and let $\Delta_{T}$ be the root system based on sub-Coxeter-Dynkin diagram corresponding to the vertices labelled by $T$. Then $V_{T} \cap \Delta=\Delta_{T}$.

Proposition 3. Let $T \subset\{1, \ldots, l\}$ be any subset. Then each orbit o of $\hat{W}_{T}$ in $\Delta$ contains a unique element $\xi \in \Delta^{+}$of the form $\Sigma n_{i} \omega_{i}$ where $n_{i} \geqslant 0$ for all $i \in T$. Furthermore, if $S=\left\{i \in T \mid n_{i}=0\right\}$ then either $|o|=\left[W_{T}: W_{S}\right]$ if $\xi \in \Delta_{T}$, or $|o|=2\left[W_{T}: W_{s}\right]$ if $\xi \notin \Delta_{T}$.

Proof (Existence). Let $o$ be an orbit of $\hat{W}_{T}$ in $\Delta$. Choose $\xi=\Sigma n_{i} \omega_{i}$ $=\Sigma c_{i} \alpha_{i} \in o$ with $h t \xi:=\Sigma c_{i}$ maximal. Since $-1_{V} \in \hat{W}_{T}, \xi \in \Delta^{+}$. Since $h t\left(r_{i} \xi\right)$ $\leqslant h t(\xi)$ for all $i \in T$, we have $\left(\xi, \alpha_{i}\right) \geqslant 0$ hence $n_{i} \geqslant 0$, if $i \in T$.
(UniQUENESS). Let $\pi: V \rightarrow V_{T}$ be the orthogonal projection onto $V_{T}=$ $\Sigma_{j \in T} \mathbf{R} \alpha_{j}$. Then $i \in T, v \in V,\left(\pi(v), \alpha_{i}\right)=\left(v, \alpha_{i}\right)$ from which it follows that $\pi$ is $W_{T}$-equivariant. Furthermore, $\pi$ is injective on any $W_{T}$-orbit in $V$ since $V_{T} \cap$ ker $\pi=0$. Thus there is a unique element on each $W_{T}$-orbit in $V$ whose projection is in the dominant chamber of $V_{T}$ under $W_{T}$. If $\mu=\Sigma u_{j} \omega_{j}$ is such an element then $\left(\mu, \alpha_{i}\right)=\left(\pi(\mu), \alpha_{i}\right) \geqslant 0$ for all $i \in T$, that is $u_{i} \geqslant 0$ for all $i \in T$.

Now let $\xi \in \Delta^{+}$. If $\hat{W}_{T} \xi=W_{T} \xi$ then $-\xi \in W_{T} \xi$ so $W_{T} \xi \cap \Delta^{-} \neq \varnothing$, from which $\xi \in \Delta_{T}^{+}$. Conversely, if $\xi \in \Delta_{T}^{+}$then $-\xi \in W_{T} \xi$ so $\hat{W}_{T} \xi=W_{T} \xi$.

Finally suppose that distinct $\xi, \xi^{\prime}$ satisfy the hypothesis of the Proposition and define the same $\hat{W}_{T}$-orbit. Then by the above $W_{T} \xi \neq W_{T} \xi^{\prime}$ so in fact $\hat{W}_{T} \xi$ $\neq W_{T} \xi, \xi \notin \Delta_{T}, W_{T} \xi \subset \Delta^{+}$and $\xi^{\prime} \in-W_{T} \xi \subset \Delta^{-}$, a contradiction. This proves the uniqueness. The statement on the orbit size is immediate.

Proposition 4 (Modified Freudenthal formula). Let $M$ be the irreducible $\left(\mathscr{V}\right.$-module of highest weight $\Lambda$ and let $\lambda \in \Omega(\Lambda)$. Let $W_{T}=\operatorname{Stab}_{W}(\lambda)$ and let $o_{1}, \ldots, o_{n}$ be the orbits of $\hat{W}_{T}$ in $\Delta \quad$ Let $\xi_{i} \in o_{i}$ be as in Proposition 3. Then equation (2) holds.

Proof. We may suppose the indexing of the orbits is taken so that $o_{1} \cup \cdots \cup o_{m}=\Delta_{T}, 0 \leqslant m \leqslant n$. For $i \leqslant m, o_{i}=W_{T} \xi_{i}$ whereas for $i>m$, $o_{i}=W_{T} \xi_{i} \cup-W_{T} \xi_{i}$ (disjoint). Begin with the form of Freudenthal's formula [3, Equation 48.2]:

$$
\begin{equation*}
c_{\Lambda} m_{\lambda}=\sum_{\alpha \in \Delta} \sum_{p=0}^{\infty}(\lambda+p \alpha, \alpha) m_{\lambda+p \alpha}+(\lambda, \lambda) m_{\lambda} \tag{6}
\end{equation*}
$$

For $w \in W_{T},(\lambda+p w \alpha, w \alpha)=(w(\lambda+p \alpha), w \alpha)=(\lambda+p \alpha, \alpha)$ and $m_{\lambda+p w \alpha}=$ $m_{\lambda+p \alpha}$, so the double sum of the right-hand side of (6) may be rewritten as

$$
\begin{align*}
& \sum_{i=1}^{m}\left|W_{T} \xi_{i}\right| \sum_{p=0}^{\infty}\left(\lambda+p \xi_{i}, \xi_{i}\right) m_{\lambda+p \xi_{i}}  \tag{7}\\
& \quad+\sum_{i=m+1}^{n}\left|W_{T} \xi_{i}\right| \sum_{p=0}^{\infty}\left\{\left(\lambda+p \xi_{i}, \xi_{i}\right) m_{\lambda+p \xi_{i}}+\left(\lambda-p \xi_{i},-\xi_{i}\right) m_{\lambda-p \xi_{i}}\right\}
\end{align*}
$$

Now for any $\alpha \in \Delta$ there is the relation [3, §48]

$$
\begin{equation*}
\sum_{p=-\infty}^{\infty}(\lambda+p \alpha, \alpha) m_{\lambda+p \alpha}=0 \tag{8}
\end{equation*}
$$

by which (7) becomes

$$
\begin{align*}
\sum_{i=1}^{m}\left|W_{T} \xi_{i}\right| & \sum_{p=0}^{\infty}\left(\lambda+p \xi_{i}, \xi_{i}\right) m_{\lambda+p \xi_{i}} \\
& +\sum_{i=m+1}^{n}\left|W_{T} \xi_{i}\right|\left\{2 \sum_{p=1}^{\infty}\left(\lambda+p \xi_{i}, \xi_{i}\right) m_{\lambda+p \xi_{i}}+\left(\lambda, \xi_{i}\right) m_{\lambda}\right\} \tag{9}
\end{align*}
$$

Collect the coefficients of $m_{\lambda}$ appearing in (9). In the first sum occur those $\xi_{i}$ which are in $\Delta_{T}$. Since $\left(\lambda, \alpha_{j}\right)=0$ for $j \in T,\left(\lambda, \xi_{i}\right)=0$ for $j=1, \ldots, m$. In the second sum we have $\Sigma_{i=m+1}^{n}\left|W_{T} \xi_{i}\right|\left(\lambda, \xi_{i}\right)=\Sigma_{\alpha \in \Delta^{+}-\Delta_{T}^{+}}(\lambda, \alpha)=$ $\Sigma_{\alpha \in \Delta^{+}}(\lambda, \alpha)=2(\lambda, \rho)$. Taking account of the ( $\left.\lambda, \lambda\right) m_{\lambda}$ appearing in (6) and $2\left|W_{T} \xi_{i}\right|=\left|o_{i}\right|$ for $i>m$, we arrive at

$$
c_{\Lambda} m_{\lambda}=\sum_{i=1}^{n}\left|o_{i}\right| \sum_{p=1}^{\infty}\left(\lambda+p \xi_{i}, \xi_{i}\right) m_{\lambda+p \xi_{i}}+(2(\lambda, \rho)+(\lambda, \lambda)) m_{\lambda}
$$

which proves the proposition.
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